

Vertical ALE coordinates and associated splitting error

Laurent Debreu, Jules Maurel-Barros

Inria, Grenoble

2D x-z tracer's advection

- In an incompressible fluid.

$$\left\{ \begin{array}{l} \frac{\partial q}{\partial t} + u \frac{\partial q}{\partial x} + w \frac{\partial q}{\partial z} = 0 \\ \frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0 \end{array} \right.$$

In generalized vertical coordinates

$$z \rightarrow s(x, z, t), h = \frac{\partial z}{\partial s}, \Omega = \frac{\partial s}{\partial t} \Big|_z + u \frac{\partial s}{\partial x} \Big|_z + w \frac{\partial s}{\partial z}$$

$$\left\{ \begin{array}{l} \frac{\partial h q}{\partial t} + \frac{\partial h u q}{\partial x} \Big|_s + \frac{\partial \Omega q}{\partial s} = 0 \\ \frac{\partial h}{\partial t} + \frac{\partial h u}{\partial x} \Big|_s + \frac{\partial \Omega}{\partial s} = 0 \end{array} \right.$$

Generalized vertical coordinates

$$\frac{\partial h}{\partial t} + \frac{\partial hu}{\partial x} \Big|_s + \frac{\partial \Omega}{\partial s} = 0$$

- Quasi Eulerian: prescribe $h(x, s, t)$ or $\frac{\partial h}{\partial t}$ and deduce Ω
- Lagrangian: set $\Omega = 0$
- ALE: Keep $\frac{\partial h}{\partial t}$ close to $-\frac{\partial hu}{\partial x}$ and try to maintain h between minimum and maximum values, with sufficient regularity

Generalized vertical coordinates

In practice, h may become small:

- Use an unconditionally stable (implicit) vertical advection scheme:

$$\frac{(hq)^{n+1} - (hq)^n}{\Delta t} + \frac{\partial huq^n}{\partial x} + \frac{\partial \Omega q^{n+1}}{\partial s} = 0 \quad \longrightarrow \quad \text{Not accurate (too dissipative)}$$

- Perform a directional splitting scheme:

$$\frac{(hq)^* - (hq)^n}{\Delta t} + \frac{\partial huq^n}{\partial x} = 0,$$

$$\frac{(hq)^{n+1} - (hq)^*}{\Delta t} + \frac{\partial \Omega q^*}{\partial s} = 0 \quad \text{with an unconditionally stable and accurate (i.e. non implicit) vertical advection scheme.}$$

↓

Flux-Form semi-lagrangian scheme advection scheme or (by the RTT), equivalent to remap q^* from h^* to h^{n+1} (since $\frac{h^{n+1} - h^*}{\Delta t} + \frac{\partial \Omega}{\partial s} = 0$)

NB: using a (high-order) semi-lagrangian scheme without splitting is not stable

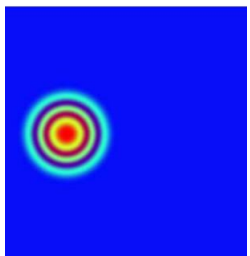
Back to a fixed vertical grid $h = \text{Constant}$

Directional splitting:

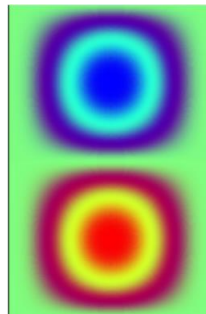
$$\frac{(q)^* - (q)^n}{\Delta t} + \frac{\partial u q^n}{\partial x} = 0,$$

$$\frac{(q)^{n+1} - (q)^*}{\Delta t} + \frac{\partial w q^*}{\partial z} = 0 \text{ (using Lax-Wendroff)}$$

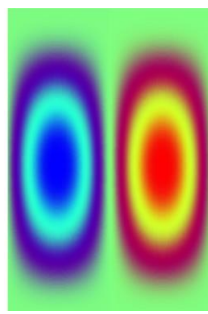
Leveque test case, $u(x, z, t) = \sin(\pi x)^2 \sin(2\pi z) \cos(\pi t/T)$



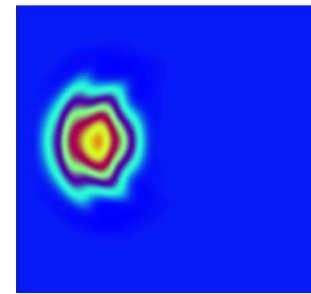
$q(t = 0)$



$u(t = 0)$



$w(t = 0)$



$q(t = 2T)$

Back to a fixed vertical grid $h = \text{Constant}$

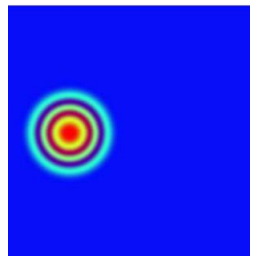
From Lie to Strang directional splitting:

$$\frac{(q)^* - (q)^n}{\Delta t/2} + \frac{\partial u q^n}{\partial x} = 0,$$

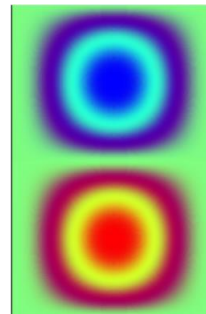
$$\frac{(q)^{**} - (q)^*}{\Delta t} + \frac{\partial w q^*}{\partial z} = 0 \text{ (using Lax-Wendroff) ,}$$

$$\frac{(q)^{n+1} - (q)^{**}}{\Delta t/2} + \frac{\partial u q^{**}}{\partial x} = 0.$$

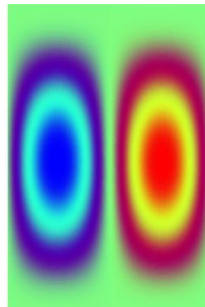
Leveque test case, $u(x, z, t) = \sin(\pi x)^2 \sin(2\pi z) \cos(\pi t/T)$



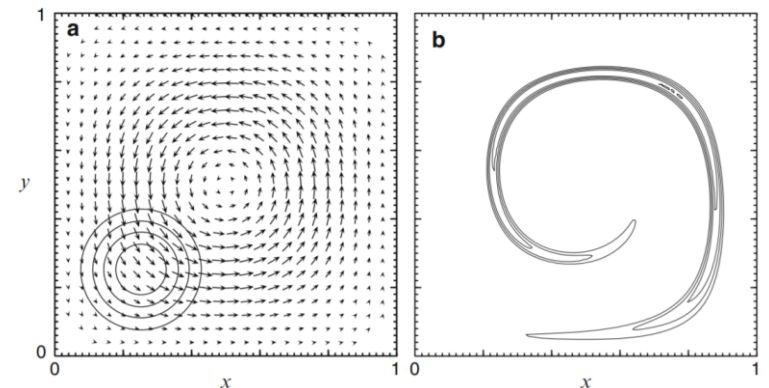
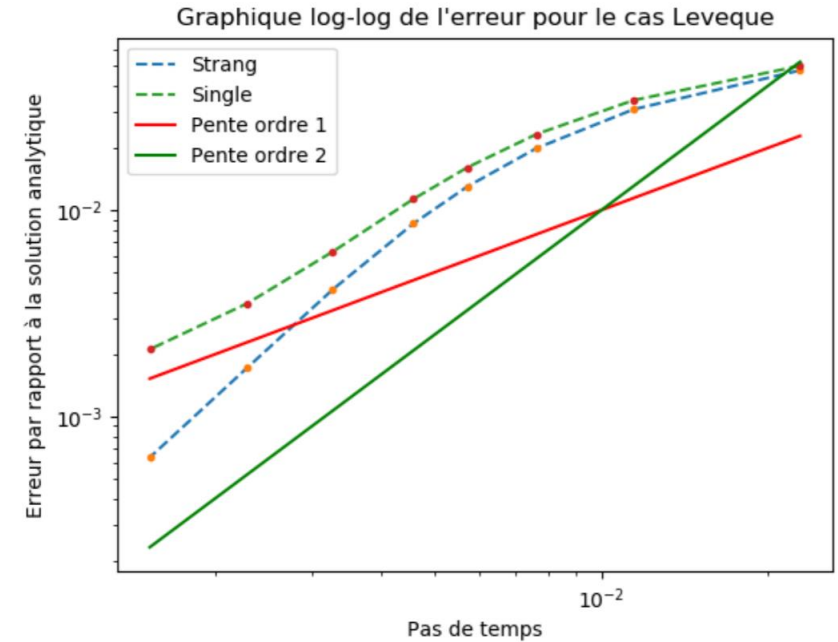
$q(t = 0)$



$u(t = 0)$



$w(t = 0)$



Back to a fixed vertical grid $h = \text{Constant}$

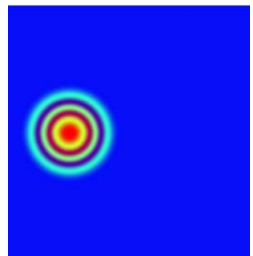
From Lie to Strang directional splitting:

$$\frac{(q)^* - (q)^n}{\Delta t/2} + \frac{\partial u q^n}{\partial x} = 0,$$

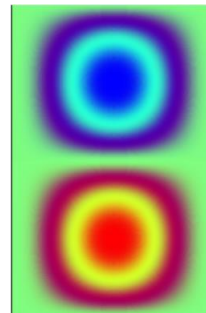
$$\frac{(q)^{**} - (q)^*}{\Delta t} + \frac{\partial w q^*}{\partial z} = 0 \text{ (using Lax-Wendroff) ,}$$

$$\frac{(q)^{n+1} - (q)^{**}}{\Delta t/2} + \frac{\partial u q^{**}}{\partial x} = 0.$$

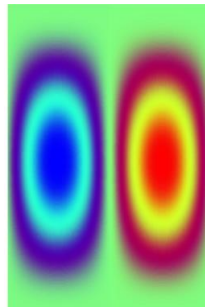
Leveque test case, $u(x, z, t) = \sin(\pi x)^2 \sin(2\pi z) \cos(\pi t/T)$



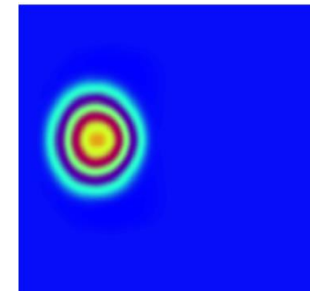
$q(t = 0)$



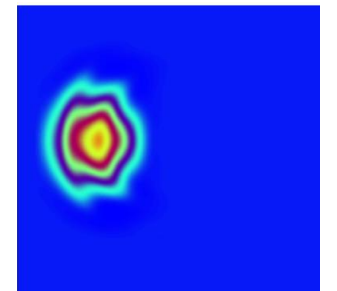
$u(t = 0)$



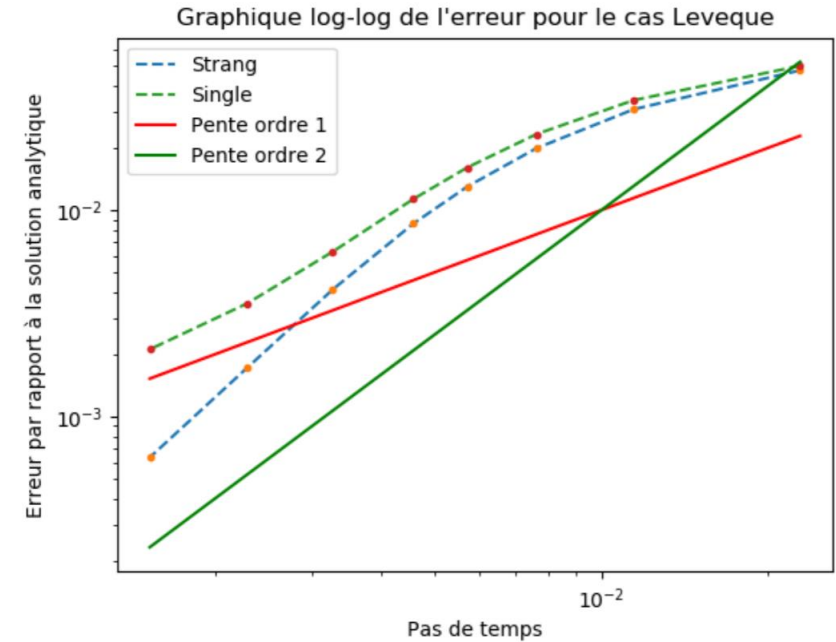
$w(t = 0)$



$q(t = 2T)$



$q(t = 2T)$



Back to a moving vertical grid ...

With a vertical Courant number exceeding 1 (at some points in space and time)

$$\frac{|\Omega|}{h} \Delta t > 1$$

$$\frac{(hq)^* - (hq)^n}{\Delta t} + \frac{\partial huq^n}{\partial x} = 0, \quad \text{RK3}$$

$$\frac{(hq)^{n+1} - (hq)^*}{\Delta t} + \frac{\partial \Omega q^*}{\partial s} = 0 \quad \text{Vertical remapping}$$

Lie and Strang versions

Back to a moving vertical grid ...

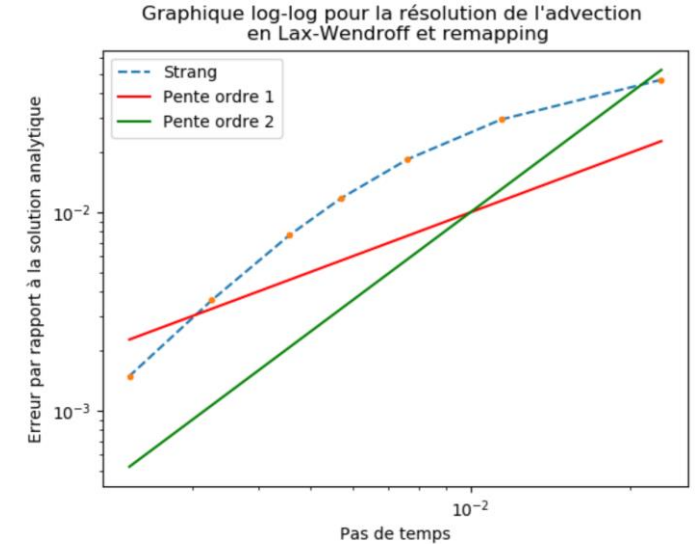
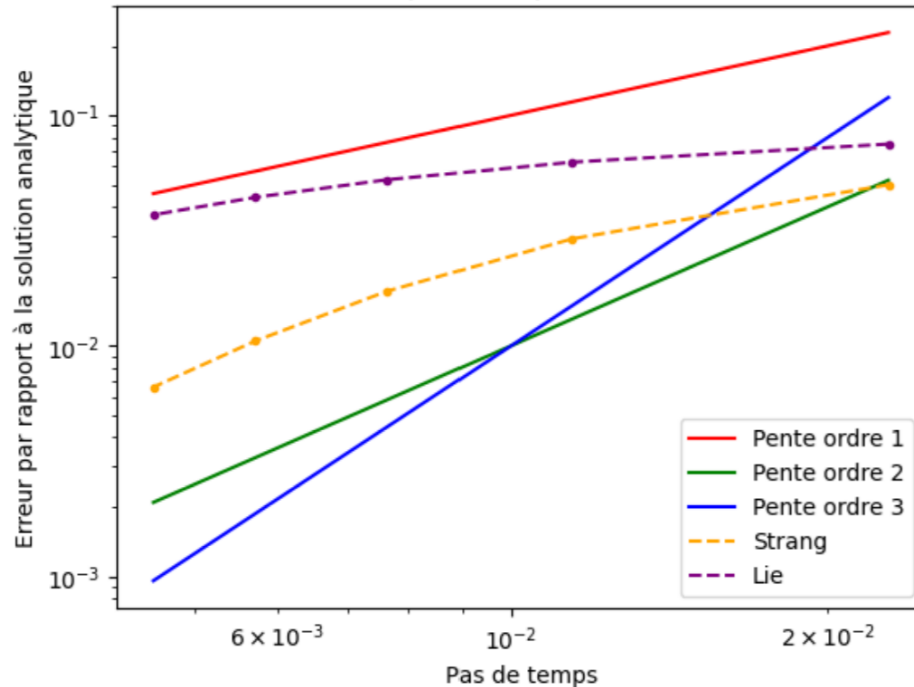
$$\frac{(hq)^* - (hq)^n}{\Delta t} + \frac{\partial huq^n}{\partial x} = 0, \quad \text{RK3}$$

$$\frac{(hq)^{n+1} - (hq)^*}{\Delta t} + \frac{\partial \Omega q^*}{\partial s} = 0 \quad \text{Vertical remapping}$$

Lie and Strang versions

Leveque test case

Graphique log-log pour la résolution de l'advection
Cas leveque en RK3 et remapping
lorsque on dépasse la CFL



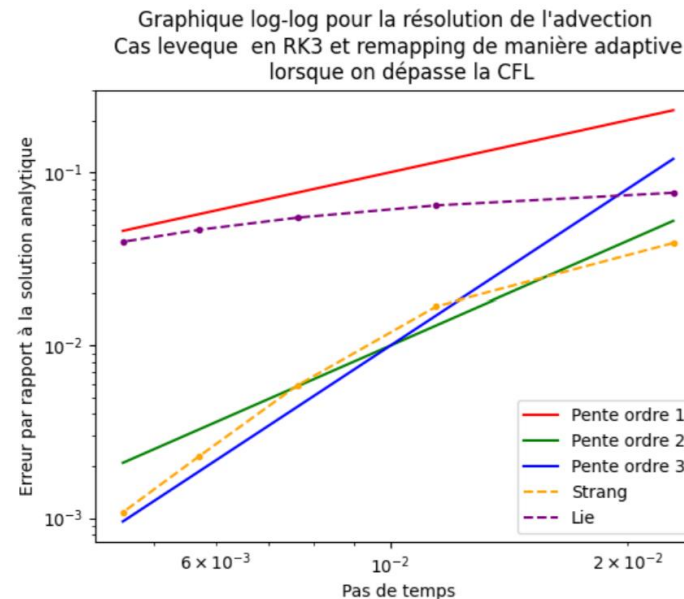
Same results for LW+remapping

A la Sasha ... $\Omega = \Omega_e + \Omega_i, \frac{|\Omega_e|}{h} < 1$

$$\frac{(hq)^* - (hq)^n}{\Delta t} + \frac{\partial huq^n}{\partial x} + \frac{\partial \Omega_e q^n}{\partial s} = 0, \quad \text{RK3}$$

$$\frac{(hq)^{n+1} - (hq)^*}{\Delta t} + \frac{\partial \Omega_i q^*}{\partial s} = 0 \quad \text{Vertical remapping}$$

Leveque test case



Strang splitting + Adaptive vertical advection allows to recover third order accuracy

What's next ?

- Perform more ocean oriented idealized advection test cases (with appropriate ALE choices)
- Tests in NEMO/CROCO are possible (the Adaptive option is already available)
- Use estimation of the splitting error (cross-derivatives terms) as an additional requirement for the design of ALE ?